

# Asymptotic stability of discrete solitons

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# The problem

The discrete nonlinear Schrödinger (DNLS) equation

$$i\dot{u}_n + \Delta_d u_n + \sigma |u_n|^2 u_n = 0, \quad n \in \mathbb{Z}^d.$$

Localized modes (time-periodic space-localized solutions) are of the form  $u_n(t) = \phi_n e^{-i\omega t}$ , where  $\omega \in \mathbb{R}$  and  $\{\phi_n\}_{n \in \mathbb{Z}^d}$  satisfies

$$(\omega + \Delta_d) \phi_n + \sigma |\phi_n|^2 \phi_n = 0, \quad n \in \mathbb{Z}^d.$$

**Main Question:** If a localized mode is orbitally stable, is it also asymptotically stable due to dispersive radiation?

- P. Kevrekidis, D. Pelinovsky, and A. Stefanov, arXiv:0810.1778
- S. Cuccagna, M. Tarulli, arXiv:0808.2024

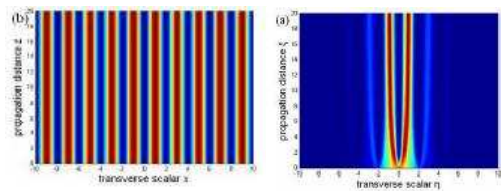
# Physical contexts

The DNLS equation arises in the modeling of **density waves in Bose–Einstein condensates** as the Gross–Pitaevskii equation

$$iu_t = -\nabla^2 u + V(x)u + \sigma|u|^2 u$$

with a bounded periodic potential  $V(x) = V(x + 2\pi)$  reduces asymptotically to the DNLS equation in a tight-binding approximation.

Another context of the DNLS equation is the **coupled waveguide arrays** in nonlinear optics and photorefractive crystals.

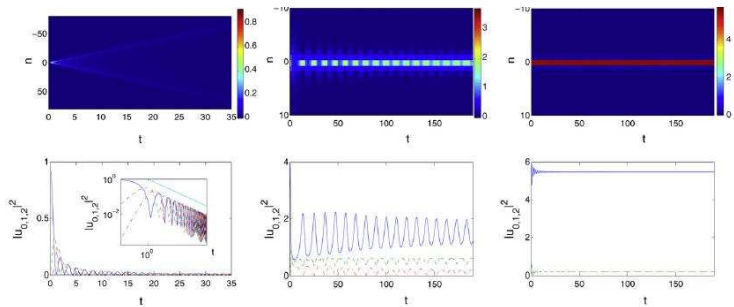


# Numerical simulations

P. Kevrekidis et al., Physics Letters A **372**, 2237 (2008)

$$(1D) \quad i\dot{u}_n + u_{n+1} - 2u_n + u_{n-1} + |u_n|^2 u_n = 0$$

$u_n(0) = A\delta_{n,0}$  with  $A = 1$  (left),  $A = 2$  (middle), and  $A = 2.5$  (right).

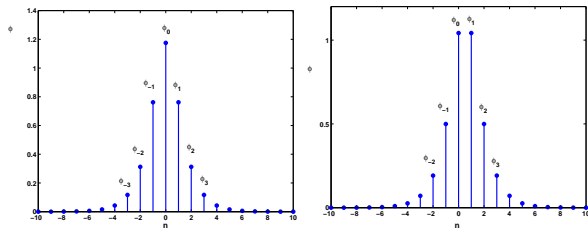


# Asymptotic stability of localized modes

Given a time-periodic space-localized solution  $\phi_n e^{-i\omega t}$  of the DNLS equation, the stability can be considered in the following three senses:

- Linearized (or spectral) stability
- Nonlinear orbital stability
- Asymptotic stability

Stability depends on  $\phi_n$ . In what follows, we consider single-humped on-site discrete solitons, which are known to be spectrally stable.



Left: on-site soliton. Right: inter-site soliton.

# More general formulation

Let us consider the 1D DNLS equation in the form

$$i\dot{u}_n = (-\Delta + V_n)u_n + \sigma|u_n|^{2p}u_n, \quad n \in \mathbb{Z},$$

where  $\sigma = \pm 1$ ,  $p \geq 1$  (an integer), and  $V \in l^\infty$ .

**Assumption:**  $V$  supports exactly one negative eigenvalue  $\omega_0 < 0$  of  $H = -\Delta + V$  with an eigenvector  $\psi_0 \in l^2$  (normalized by  $\|\psi_0\|_{l^2} = 1$ ).

For instance, if  $V_n = -\delta_{n,0}$ , the assumption is satisfied with

$$(\psi_0)_n = e^{-\kappa|n|}, \quad n \in \mathbb{Z},$$

where  $\kappa = \operatorname{arcsinh}(2^{-1})$  and  $\omega_0 = 2 - \sqrt{5} < 0$ .

# Global well-posedness

**Global well-posedness:** For any  $\mathbf{u}_0 \in l^2_\sigma$ ,  $\sigma \geq 0$ , there exists a unique solution  $\mathbf{u}(t) \in C^1(\mathbb{R}_+, l^2_\sigma)$  s.t.  $\mathbf{u}(0) = \mathbf{u}_0$  and  $\mathbf{u}(t)$  depends continuously on  $\mathbf{u}_0$ .

Local existence follows from the Picard iterations applied to

$$u_n(t) = u_n(0) - i \int_0^t [(-\Delta + V_n)u_n(t') + \sigma|u_n(t')|^{2p}u_n(t')] dt'$$

in space  $C([0, T], l^2_\sigma)$ . To show that  $T = \infty$ , we can use the balance equation

$$i \frac{d}{dt} |u_n|^2 = u_n(\bar{u}_{n+1} + \bar{u}_{n-1}) - \bar{u}_n(u_{n+1} + u_{n-1}),$$

so that

$$\|\mathbf{u}(t)\|_{l^2_\sigma}^2 \leq \|\mathbf{u}(0)\|_{l^2_\sigma}^2 + C \int_0^t \|\mathbf{u}(t')\|_{l^2_\sigma}^2 dt'.$$

By Gronwall's inequality,  $\|\mathbf{u}(t)\|_{l^2_\sigma}^2$  is bounded and continuous for any  $t \in \mathbb{R}_+$ .

# Local bifurcation of localized modes

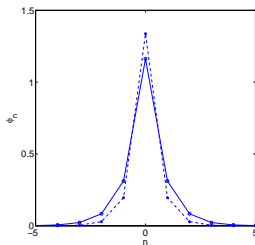
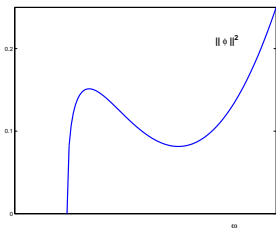
**Local bifurcation:** Let  $\epsilon := \omega - \omega_0$  and  $\sigma = +1$ . For any  $\epsilon \in (0, \epsilon_0)$ , where  $\epsilon_0 > 0$  is small, there exists a solution  $\phi(\omega) \in C([\omega_0, \omega_0 + \epsilon_0], \ell^2)$  of

$$(-\Delta + V_n)\phi_n(\omega) + \sigma\phi_n^{2p+1}(\omega) = \omega\phi_n(\omega), \quad n \in \mathbb{Z},$$

satisfying

$$\left\| \phi(\omega) - \frac{\epsilon^{\frac{1}{2p}} \psi_0}{\|\psi_0\|_{\ell^{2p+2}}^{1+\frac{1}{p}}} \right\|_{\ell^2} \leq C\epsilon^{1+\frac{1}{2p}}.$$

Moreover,  $\phi(\omega)$  decays exponentially to zero as  $|n| \rightarrow \infty$ .





**Theorem** [Weinstein, 1999]: Let  $\sigma = -1$  and  $V \equiv 0$ . There exists a global minimizer of energy

$$E = \sum_{n \in \mathbb{Z}} |u_{n+1} - u_n|^2 - \frac{1}{p+1} |u_n|^{2p+2}$$

under a fixed  $N = \|u\|_p^2$  for any  $p \geq 1$ . If  $p < 2$ , it exists for any  $N > 0$ , whereas for  $p \geq 2$ , there is a threshold  $N_0 > 0$  so that it exists for  $N \geq N_0$ .

If  $\mathbf{u}(0) \approx \phi(\omega(0))$ , then  $\mathbf{u}(t)$  remains near  $\phi(\omega(t))$  for all  $t \in \mathbb{R}_+$ . However, the question is if there exists  $\omega_\infty$  so that  $\mathbf{u}(t) \rightarrow \phi(\omega_\infty)$  as  $t \rightarrow \infty$ .

# Main result

**Theorem** [Kevrekidis, P., Stefanov, 2008]: Let  $\sigma = +1$  and  $p \geq 3$ . Fix  $\epsilon > 0$  and  $\delta > 0$  be small and assume that  $\omega(0) = \omega_0 + \epsilon$  and

$$\|\mathbf{u}(0) - \phi(\omega_0 + \epsilon)\|_{l^2} \leq \delta \epsilon^{\frac{1}{2p}}.$$

Under some assumptions on  $V$ , there exist  $\omega_\infty \in (\omega_0, \omega_0 + \epsilon_0)$ ,  $(\omega, \theta) \in C^1(\mathbb{R}_+)$ , and  $\mathbf{y}(t) = \mathbf{u}(t) - e^{-i\theta(t)}\phi(\omega(t)) \in C^1(\mathbb{R}_+, l^2) \cap L^6(\mathbb{R}_+, l^\infty)$  such that  $\mathbf{u}(t)$  solves the DNLS equation and

$$\lim_{t \rightarrow \infty} \omega(t) = \omega_\infty, \quad \lim_{t \rightarrow \infty} \|\mathbf{u}(t) - e^{-i\theta(t)}\phi(\omega(t))\|_{l^\infty} = 0.$$

**Remark:** A similar result applies in the focusing case  $\sigma = -1$  with the local bifurcation to  $\omega < \omega_0$ .

Pioneer works on continuous NLS equations are by **Soffer, Weinstein** (1990, 1992, 2004), **Pillet, Wayne** (1997), **Yao, Tsai** (2002), **Mizumachi** (2008).

# Decomposition of the solution

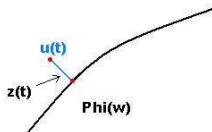
Let

$$\mathbf{u}(t) = e^{-i\theta(t)} (\phi(\omega(t)) + \mathbf{z}(t)),$$

for some  $(\omega, \theta) \in \mathbb{R}^2$ . Then,  $\mathbf{z}(t) \in C^1(\mathbb{R}_+, l_\sigma^2)$  solves

$$i\dot{\mathbf{z}} = (H - \omega)\mathbf{z} - (\dot{\theta} - \omega)(\phi(\omega) + \mathbf{z}) - i\dot{\omega}\partial_\omega\phi(\omega) + \mathbf{N}(\phi(\omega) + \mathbf{z}) - \mathbf{N}(\phi(\omega)),$$

where  $H = -\Delta + V$  and  $[\mathbf{N}(\psi)]_n = \sigma|\psi_n|^{2p}\psi_n$ .



**Question:** How to ensure that the decomposition is unique?

# Double null space

Linearized time evolution for  $\mathbf{z}(t) = \mathbf{v}(t) + i\mathbf{w}(t)$  is defined by the non-self-adjoint eigenvalue problem

$$L_+ \mathbf{v} = -\lambda \mathbf{w}, \quad L_- \mathbf{w} = \lambda \mathbf{v},$$

where

$$L_- = H - \omega + W, \quad L_+ = H - \omega + (2p + 1)W,$$

where  $W_n = \sigma \phi_n^{2p}(\omega)$ .

There exists a double zero eigenvalue with a one-dimensional kernel, isolated from the rest of the spectrum. The generalized kernel is spanned by vectors  $(\mathbf{0}, \phi(\omega)), (-\partial_\omega \phi(\omega), \mathbf{0}) \in \ell^2$  satisfying

$$L_- \phi(\omega) = \mathbf{0}, \quad L_+ \partial_\omega \phi(\omega) = \phi(\omega).$$

$(\mathbf{v}, \mathbf{w}) \in \ell^2$  is symplectically orthogonal to the double subspace of the generalized kernel under the conditions

$$\langle \mathbf{v}, \phi(\omega) \rangle = 0, \quad \langle \mathbf{w}, \partial_\omega \phi(\omega) \rangle = 0,$$

where  $\langle \mathbf{u}, \mathbf{v} \rangle := \sum_{n \in \mathbb{Z}} u_n \bar{v}_n$ .

# Symplectic orthogonality

Let us require that

$$\langle \operatorname{Re} \mathbf{z}(t), \psi_1 \rangle = \langle \operatorname{Im} \mathbf{z}(t), \psi_2 \rangle = 0,$$

where

$$\psi_1 = \frac{\phi(\omega)}{\|\phi(\omega)\|_{\mathcal{L}^2}}, \quad \psi_2 = \frac{\partial_\omega \phi(\omega)}{\|\partial_\omega \phi(\omega)\|_{\mathcal{L}^2}}.$$

**Unique decomposition:** Fix  $\epsilon \in (0, \epsilon_0)$ . There exists  $\delta > 0$  and  $T > 0$  such that any  $\mathbf{u}(t) \in C^1([0, T], \mathcal{L}^2)$  satisfying

$$\|\mathbf{u}(t) - \phi(\omega_0 + \epsilon)\|_{\mathcal{L}^2} \leq \delta \epsilon^{\frac{1}{2p}}, \quad t \in [0, T],$$

can be uniquely decomposed by

$$\mathbf{u}(t) = e^{-i\theta(t)} (\phi(\omega(t)) + \mathbf{z}(t)),$$

where  $(\omega, \theta) \in C^1([0, T], \mathbb{R}^2)$  and  $\mathbf{z}(t) \in C^1([0, T], \mathcal{L}^2)$  satisfies the symplectic orthogonality conditions. Moreover, there exists  $C > 0$  such that

$$|\omega(t) - \omega_0 - \epsilon| \leq C\delta\epsilon, \quad |\theta(t)| \leq C\delta, \quad \|\mathbf{z}(t)\|_{\mathcal{L}^2} \leq C\delta\epsilon^{\frac{1}{2p}}, \quad t \in [0, T].$$

# Projections

The time-evolution of  $(\omega, \theta)$  satisfies the system

$$\mathbf{A}(\omega, \mathbf{z}) \begin{bmatrix} \dot{\omega} \\ \dot{\theta} - \omega \end{bmatrix} = \mathbf{f}(\omega, \mathbf{z}),$$

where

$$\mathbf{A}(\omega, \mathbf{z}) = \begin{bmatrix} \langle \partial_{\omega} \phi(\omega), \psi_1 \rangle - \langle \operatorname{Re} \mathbf{z}, \partial_{\omega} \psi_1 \rangle & \langle \operatorname{Im} \mathbf{z}, \psi_1 \rangle \\ \langle \operatorname{Im} \mathbf{z}, \partial_{\omega} \psi_2 \rangle & \langle \phi(\omega) + \operatorname{Re} \mathbf{z}, \psi_2 \rangle \end{bmatrix}$$

and

$$\mathbf{f}(\omega, \mathbf{z}) = \begin{bmatrix} \langle \operatorname{Im} \mathbf{N}(\phi + \mathbf{z}) - \mathbf{W} \mathbf{z}, \psi_1 \rangle \\ \langle \operatorname{Re} \mathbf{N}(\phi + \mathbf{z}) - \mathbf{N}(\phi) - (2\rho + 1) \mathbf{W} \mathbf{z}, \psi_2 \rangle \end{bmatrix}.$$

In addition, we recall the time evolution of  $\mathbf{z}(t)$  from

$$i\dot{\mathbf{z}} = (H - \omega)\mathbf{z} - (\dot{\theta} - \omega)(\phi(\omega) + \mathbf{z}) - i\dot{\omega}\partial_{\omega}\phi(\omega) + \mathbf{N}(\phi(\omega) + \mathbf{z}) - \mathbf{N}(\phi(\omega)).$$

# Dispersive decay estimates

**Assumption:** Let  $V \in L^1_{2\sigma}$  for a fixed  $\sigma > \frac{5}{2}$  and let  $V$  is generic in the sense that no solution  $\psi_0$  of equation  $H\psi_0 = 0$  exists in  $L^2_{-\sigma}$  for  $\frac{1}{2} < \sigma \leq \frac{3}{2}$ .

**Pointwise dispersive decay estimates:** There exists a constant  $C > 0$  depending on  $V$  such that

$$\begin{aligned} \left\| \langle n \rangle^{-\sigma} e^{-itH} P_{a.c.}(H) \mathbf{f} \right\|_{l^2_n} &\leq C(1+t)^{-3/2} \|\langle n \rangle^\sigma \mathbf{f}\|_{l^2_n}, \\ \left\| e^{-itH} P_{a.c.}(H) \mathbf{f} \right\|_{l^\infty_n} &\leq C(1+t)^{-1/3} \|\mathbf{f}\|_{l^1_n}, \end{aligned}$$

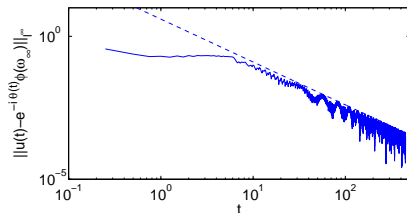
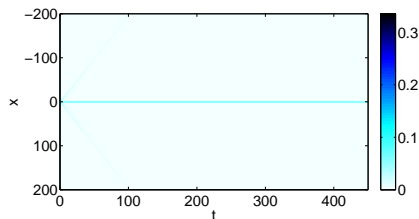
for all  $t \in \mathbb{R}_+$ .

**Discrete Strichartz estimates:** There exists a constant  $C > 0$  such that

$$\begin{aligned} \left\| e^{-itH} P_{a.c.}(H) \mathbf{f} \right\|_{L_t^6 l_n^\infty \cap L_t^\infty l_n^2} &\leq C \|\mathbf{f}\|_{l_n^2}, \\ \left\| \int_0^t e^{-i(t-s)H} P_{a.c.}(H) \mathbf{g}(s) ds \right\|_{L_t^6 l_n^\infty \cap L_t^\infty l_n^2} &\leq C \|\mathbf{g}\|_{L_t^1 l_n^2}. \end{aligned}$$

# Numerical results

For any  $p = 1, 2, 3$ , it was found that  $\|\mathbf{y}(t)\|_{l^\infty} = O(t^{-3/2})$  as  $t \rightarrow \infty$ .



## Further results:

- Cuccagna' 09: periodic oscillations of discrete solitons with  $V$  supporting two eigenvalues
- Stefanov' 09: pushing analysis to  $p = 2$